Fluorescence enhancement and structural development of sol-gel derived Er³⁺-doped SiO₂ by yttrium codoping

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 $Er^{3+}-Y^{3+}$ codoped SiO_2 powdered bulks were prepared by a sol-gel process. The effect of Y^{3+} codoping on the fluorescence properties and structural development of Er^{3+} -doped SiO_2 is investigated. The maximum $\sim 1.54~\mu m$ photoluminescence (PL) intensity occurs in the sample with Er^{3+} (10 mol%)– Y^{3+} (50 mol%) codoped and annealed at 985 °C. This can be attributed to the competition between the content of hydroxy groups and Er site symmetry. Improvement of optimum PL properties due to Y^{3+} codoping by a factor of ~ 20 for intensity and 1.8 for the full width at half maximum (57 nm) was obtained in comparison with the Er^{3+} -doped SiO_2 system. Extended X-ray absorption fine structure analysis shows that the local chemical environment of Er^{3+} ions in the $Er^{3+}-Y^{3+}$ codoped SiO_2 is similar to that in Er_2O_3 . The average spatial distance between Er^{3+} ions is enlarged due to a partial substitution of Y^{3+} for Er^{3+} ions in the Er_2O_3 -like local structure, which results in a reduction of the concentration quenching effect.

1. Introduction

Rare-earth doped materials play very important roles in optoelectronics technology because they have good performance in the application of lasers and optical amplifiers. $^{1-3}$ Among these rare-earth elements, the $\sim\!1.54~\mu m$ photoluminescence (PL) properties of $\rm Er^{3+}$ ions are of particular interest in fiber optical communication systems because it matches the lowest signal attenuation in silica-based optical fibers. $^{4.5}$

The $\sim 1.54 \,\mu m$ PL efficiency was strongly influenced by some factors such as the local symmetry of Er3+ sites in the host matrix, the concentration quenching effect, and the content of hydroxy impurities (OH⁻) in the sol-gel derived materials.^{6,7} In the case of free Er³⁺ ions, the $\sim 1.54 \mu m$ optical emission of internal 4f–4f transitions (between the ${}^4I_{13/2}$ first excited and the ${}^4I_{15/2}$ ground state) is electric dipole forbidden. If the symmetry of the local crystal field around the Er³⁺ lattice sites in the host matrix is distorted, the parity forbidden intra-4f transition will be allowed.^{8,9} Additionally, for a high Er³⁺ doping, the Er³⁺ ions tend to form clusters and the mean interaction distance between Er3+ ions becomes small, which results in substantial losses of excitation energy through the ion-ion interaction mechanism (e.g., co-operative up-conversion cross-relaxation and resonant energy migration) between the two nearby Er³⁺ ions. In other words, the chemical environment of Er³⁺ ions (i.e. the site symmetry and clustering of Er³⁺ ions) in the host matrix significantly affects the luminescence efficiency. Therefore, the addition of specific heteroatoms into the host matrix to reduce the site symmetry and clustering of Er³⁺ ions will be the most efficient method to promote the $\sim 1.54 \mu m$ PL intensity. For instance, it is well known that the solubility of Er3+ ions in silicate hosts is very low. Many researchers have tried to codope Al³⁺ or P³⁺ ions with Er³⁺ ions into a silicate host in order to provide enough non-network oxygen species and hence decrease the concentration quenching effect. ^{10–13} The addition of co-dopants Al³⁺ or P³⁺ ions has focused on the modification of the host matrix to improve the Er solubility. However, the structural role of Er in the host matrix has been neglected.

 Y^{3+}/Er^{3+} ions have similar ionic radii ($Y^{3+} = 0.0892$, and

 $\rm Er^{3+}=0.0881$ nm) and $\rm Y_2O_3/Er_2O_3$ have nearly the same crystal structure and lattice constant. Therefore, we have tried to codope $\rm Y^{3+}$ ions into the $\rm Er^{3+}$ -doped $\rm SiO_2$ network by the sol–gel method in the present work. The role of the $\rm Y^{3+}$ codopant in the phase development and PL properties of $\rm Er^{3+}$ -doped silica materials will be investigated. Additionally, the extended X-ray absorption fine structure (EXAFS) technique will be used to analyze the influence of $\rm Y^{3+}$ ions on the local chemical environment around the $\rm Er^{3+}$ ions in an $\rm Er^{3+}$ - $\rm Y^{3+}$ codoped $\rm SiO_2$ system.

2. Experimental

2.1. Sample preparation

TEOS [tetraethoxysilane, Si(OC₂H₅)₄, Merck, 99.5%] was used as the silicon alkoxide precursor for all of the Er³⁺-doped samples. Erbium nitrate pentahydrate [Er(NO₃)₃·5H₂O, Alfa, 99%], and yttrium nitrate pentahydrate [Y(NO₃)₃·5H₂O, Alfa, 99%] were used as the sources of Er^{3+} and Y^{3+} ions, respectively. TEOS was first added to the mixture solution of deionized water and ethanol (Merck, 99.9%) (the molar ratio of TEOS: $H_2O: C_2H_5OH = 1:5:10$), followed by stirring for 30 min. Then the erbium solution (erbium nitrate: H_2O = 1:10) and yttrium solution (yttrium nitrate: $H_2O = 1:10$) were spontaneously dropped into the initial TEOS solution. This sol solution was vigorously stirred at room temperature for 10 h in order to ensure a sufficient degree of hydrolysis and polycondensation. Molar ratios of Er³⁺ and Y³⁺ (relative to Si⁴⁺) ranging from 0.5 to 10% and 10 to 50% were used, respectively.

The final sol solutions were aged at room temperature for 6 days and gelled at 50 °C in about 2 weeks. Gels were heattreated at 100 °C in air for 6 days and were then pyrolyzed at 700 °C for 1 h in a dry oxygen atmosphere at a heating rate of 2 °C h⁻¹ to remove organic species. Pyrolyzed gel powder was die-pressed to be a 1 cm-diameter pellet and then annealed at 800–1400 °C for 1 h in a dry oxygen atmosphere at a heating rate of 10 °C min⁻¹.

2.2. Characterization measurements

The phase structure of the samples was determined by an X-ray diffractometer (MAC Science, M18X) using Cu-K α radiation. Fourier transform infrared (FTIR) spectra were recorded on a Perkin–Elmer 580 spectrometer. A transmission electron microscope (TEM, JEOL-2000FX) equipped with EDX was used to observe and analyze the microstructure. The fluorescence spectra were excited by a 980 nm laser diode with a power of 50 mW inclined 45° to irradiate the samples and recorded normally from the samples using a spectrophotometer equipped with a liquid N₂-cooled Ge detector (NCSC).

Erbium $L_{\rm III}$ -edge X-ray absorption spectra were recorded on Wiggler beamline S-05B at the Synchrotron Radiation Research Center (SRRC), Hsinchu, Taiwan. The electron storage ring was operated at an energy of 1.3 GeV and a current of 80–200 mA. A Si (111) double-crystal monochromator with a 0.5 mm entrance slit was used for energy scanning. The energy resolution, $\Delta E/E$, was about 1.9×10^{-4} . Measurements were performed at room temperature in fluorescence mode and the sample was positioned at 45° to the incident X-ray beam. A polycrystalline $\rm Er_2O_3$ powder (Cerac, 99.9%) was used as a reference standard.

3. Results

The X-ray diffraction (XRD) patterns in Fig. 1 illustrate the effect of the annealing temperatures on the structural evolution of pure SiO₂, Er³⁺-doped SiO₂ and Y³⁺-doped SiO₂, where the silicate host codoped with Er³⁺ (20 mol%)–Y³⁺ (0 mol%) and annealed at 1100 °C is denoted as "Er20 Y0 1200 °C". For annealing temperatures below 1300 °C, the pure SiO₂ host exhibits an amorphous structure. The polymorphic SiO₂ is detectable when the annealing temperature exceeds 1400 °C. When 10 mol% Er³⁺ ions were added into the SiO₂ matrix, it was found that phase crystallization occurs at 1100 °C. With increasing Er³⁺ concentration up to 20 mol%, the well-defined crystalline phase can be clearly identified as the Erpyrosilicate phase $(Er_2Si_2O_7)$. ¹⁴ On the other hand, Y^{3+} -doped SiO₂ exhibits similar crystallization behavior to Er³⁺-doped SiO_2 and also forms the polymorphic Y-pyrosilicate phase $(Y_2Si_2O_7)$ at 1100 °C. ¹⁵ Notedly, both $Er_2Si_2O_7$ and $Y_2Si_2O_7$ phases have almost the same XRD patterns, which means that both Er- and Y-pyrosilicates have the same crystal structure and nearly the same lattice constant.

The effect of yttrium concentration on the phase evolution of the $\mathrm{Er^{3^+}}$ (5 mol%)-doped $\mathrm{SiO_2}$ at different annealing temperatures is shown in Fig. 2. When annealed at 985 °C/1 h, even though as high as 50 mol% $\mathrm{Y^{3^+}}$ was added into the $\mathrm{Er^{3^+}}$ -doped $\mathrm{SiO_2}$, the XRD pattern shows an amorphous

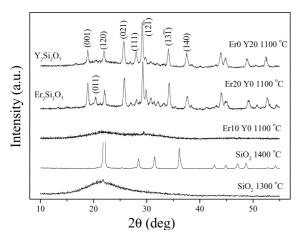


Fig. 1 X-Ray diffraction patterns of SiO_2 and Er^{3+} (or Y^{3+})-doped SiO_2 samples annealed at different temperatures for 1 h.

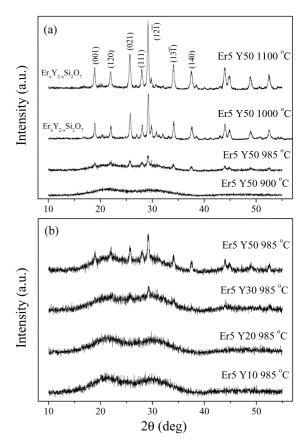


Fig. 2 X-Ray diffraction patterns of Er^{3+} (5 mol%)– Y^{3+} codoped SiO_2 samples (a) annealed at temperatures of 900–1000 °C/1 h, and (b) with different Y^{3+} concentration and annealed at 985 °C/1 h.

structure. The relationship between the total $Er^{3+} + Y^{3+}$ codoping concentration and annealing temperatures for the crystallization behavior of $Er^{3+} - Y^{3+}$ codoped SiO_2 samples is further summarized in Fig. 3. For example, when the total concentration of Er^{3+} and Y^{3+} ions is more than 30 mol%, the crystalline phase $(Er_xY_{2-x}Si_2O_7)$ will be generated for all of the samples annealed at 985 °C.

Fig. 4 shows the FTIR spectra of the ${\rm Er^{3+}-Y^{3+}}$ codoped ${\rm SiO_2}$ samples annealed at 800 and 1000 °C for 1 h. The FTIR spectrum of ${\rm Er^{3+}}$ (5 mol%)– ${\rm Y^{3+}}$ (10 mol%) codoped ${\rm SiO_2}$ annealed at 800 °C is found to be similar to that of the pure silicate. The bands around 1079 cm $^{-1}$ and 796 cm $^{-1}$ correspond to the Si–O–Si antisymmetrical and symmetrical stretching vibrations, respectively. 16,17,18 The band around 455 cm $^{-1}$ is due to the Si–O–Si and O–Si–O bending modes. $^{18-20}$ However, with increasing ${\rm Y^{3+}}$ concentration (from 30 to

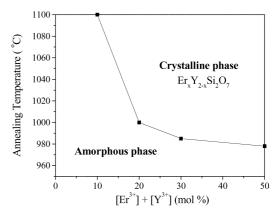


Fig. 3 Dependence of crystallization behavior of ${\rm Er}^{3+} - {\rm Y}^{3+}$ codoped ${\rm SiO}_2$ samples on total ${\rm Er}^{3+} + {\rm Y}^{3+}$ codoping concentration and annealing temperatures.

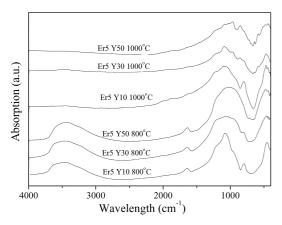


Fig. 4 FTIR absorption spectra of ${\rm Er^{3+}-Y^{3+}}$ codoped ${\rm SiO_2}$ samples annealed at different temperatures for 1 h.

50 mol%), the band around 1079 cm $^{-1}$ becomes broad and the band around 796 cm $^{-1}$ gradually weakens and disappears. This reveals that the network structure of short-range order tetrahedral SiO₄ was destroyed and became more disordered. In addition, all of the samples exhibit a broad absorption in the 3000–3450 cm $^{-1}$ region (the O–H stretching vibration) and a small band at ~ 1620 cm $^{-1}$ (the molecular H–O–H bending mode). This indicates that there exists a variety of hydroxy groups such as isolated Si–OH groups, pairs of hydrogen bonded Si–OH and physically adsorbed water molecules in the silica matrix. 18,21,22

As the sample was annealed at 1000 °C, some newly-formed peaks at 585, 851, 966 and 1084 cm⁻¹ in the FTIR spectra were observed for the samples with a Y^{3+} concentration above 30 mol% that could be related to the formation of the crystalline pyrosilicate Er_xY_{2-x}Si₂O₇ phase. However, the assignment of these peaks is not the focus of this paper. Additionally, the absorption bands around 3000-3450 and 1620 cm⁻¹ have obviously disappeared. Since the content of hydroxy quenching centers is sensitive to the annealing temperature, the FTIR spectra in the range of 4000-3000 cm⁻¹ (where the intensity of absorption bands represents the content of hydroxy groups) were performed for the Er³⁺ (5 mol%)– Y³⁺ (50 mol%) codoped SiO₂ samples annealed at 800-1000 °C. Fig. 5 illustrates that the content of hydroxy groups is rapidly reduced with an increase of annealing temperature from 970 to 1000 °C.

Fig. 6 illustrates the $\sim 1.54~\mu m$ PL spectra of the Er $^{3+}$ (1 and 5 mol%)–Y $^{3+}$ (0–50 mol%) codoped SiO $_2$ samples annealed at 985 °C for 1 h. These spectra exhibit a broad PL emission that consists of a $\sim 1.538~\mu m$ main peak and some broad shoulders. Moreover, the spectral bandwidths also become enlarged with

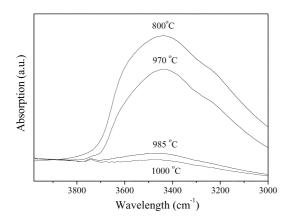


Fig. 5 FTIR absorption spectra between 4000 and 3000 cm^{$^{-1}$} showing the effect of annealing temperature on the hydroxy content of Er^{$^{3+}$} (5 mol%)–Y^{$^{3+}$} (50 mol%) codoped SiO₂ samples.

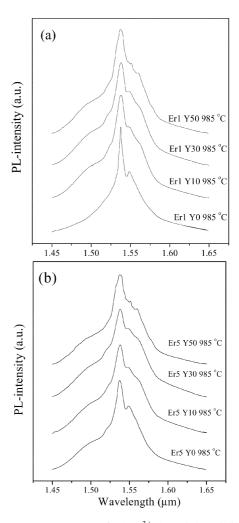


Fig. 6 \sim 1.54 μ m PL spectra of (a) Er³⁺ (1 mol%) and (b) Er³⁺ (5 mol%)–doped SiO₂ samples with different Y³⁺ concentration codoped and annealed at 950 °C for 1 h. All of the spectra are normalized on the same basis of intensity for comparison among spectral features.

more Y^{3+} codopant being added. Fig. 7 shows the annealing temperature dependence of the $\sim 1.54~\mu m$ PL spectra observed from the Er³⁺ (5 mol%)–Y³⁺ (50 mol%) codoped SiO₂ samples. The broad PL spectra will split into many sharp lines when the crystalline Er_xY_{2-x}Si₂O₇ phase is generated in the host matrix at 1000 °C

The effect of annealing temperature on the PL intensity of the ${\rm Er}^{3+}$ (1–10 mol%)– ${\rm Y}^{3+}$ (50 mol%) codoped ${\rm SiO}_2$ samples is further illustrated in Fig. 8. The results reveal that the PL intensities were enhanced with increasing annealing temperatures from 800 to 985 °C, while it decreased above 1000 °C. The dependence of ${\rm Er}^{3+}$ and ${\rm Y}^{3+}$ concentration on the PL intensity of ${\rm Er}^{3+}$ – ${\rm Y}^{3+}$ codoped ${\rm SiO}_2$ samples annealed at 985 °C for 1 h is schematically summarized in Fig. 9. The PL intensity of the ${\rm Er}^{3+}$ (5 and 10 mol%)– ${\rm Y}^{3+}$ codoped ${\rm SiO}_2$ samples increases remarkably with the increase of the ${\rm Y}^{3+}$ codoping level. However, for the ${\rm Er}^{3+}$ (1 mol%)– ${\rm Y}^{3+}$ codoped ${\rm SiO}_2$ samples, the PL intensities are slightly reduced when the ${\rm Y}^{3+}$ codoping level exceeds 30 mol%. Notably, the PL intensity of the ${\rm Er}^{3+}$ (5 and 10 mol%)-doped ${\rm SiO}_2$ sample can be increased by a factor of almost 20 in the presence of 50 mol% ${\rm Y}^{3+}$ codopant, demonstrating that ${\rm Y}^{3+}$ codoping is a very efficient method for enhancing the PL intensity in the ${\rm Er}^{3+}$ -doped ${\rm SiO}_2$ system.

Fig. 10 shows the pseudo-radial distribution functions obtained from the k^3 -weighted Fourier transforms of the ${\rm Er}^{3+}{-}{\rm Y}^{3+}$ codoped SiO₂ samples annealed at 900–1000 °C for 1 h. Qualitative observation reveals that the first-neighbor

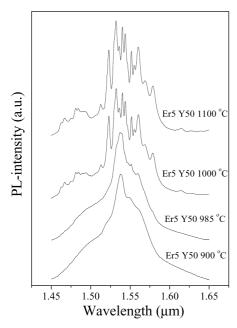


Fig. 7 \sim 1.54 µm PL spectra of Er³⁺ (5 mol%)–Y³⁺ (50 mol%) codoped SiO₂ samples annealed at 900–1100 °C for 1 h. All of the spectra are normalized on the same basis of intensity for comparison among spectral features.

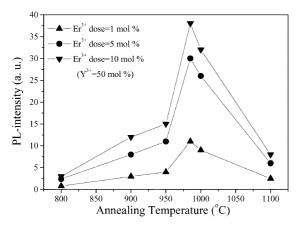


Fig. 8 PL intensity of $\rm Er^{3+}-Y^{3+}$ (50 mol%) codoped $\rm SiO_2$ samples as functions of $\rm Er^{3+}$ content and annealing temperatures.

distance of Er^{3+} at 900 °C is similar to that of the crystalline Er_2O_3 , indicating that the first neighbors of Er^{3+} ions are primarily composed of O atoms. The second nearest neighbor distance of Er^{3+} is close to the distance of the second nearest

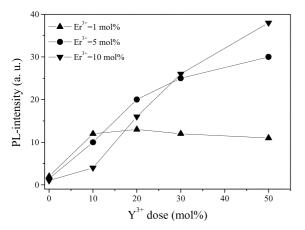


Fig. 9 Dependence of Er^{3+} and Y^{3+} concentration on the PL intensity of Er^{3+} – Y^{3+} codoped SiO_2 samples annealed at 985 °C for 1 h.

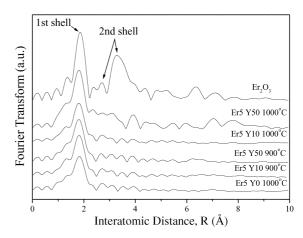


Fig. 10 Fourier transform partial radial distribution function for the Er³⁺–Y³⁺ codoped SiO₂ samples annealed at 900–1100 °C for 1 h.

neighbor in Er₂O₃. Therefore, the local chemical environment of Er³⁺ ions in the amorphous Er³⁺–Y³⁺ codoped SiO₂ samples is an Er₂O₃-like environment. At 1000 °C, the EXAFS curves of the Er³⁺ (5 mol%)–Y³⁺ (50 mol%) codoped SiO₂ samples illustrate that the first shell is like that of Er₂O₃ but its outer shells are obviously different from those of amorphous samples because the coordination number of the Er³⁺ ion in the well-crystalline Er_xY_{2-x}Si₂O₇ phase is still 6-fold. ^{23,24}

4. Discussion

4.1. Influence of \mathbf{Y}^{3+} content and annealing temperature on the PL properties

When the Er^{3+} (1–10 mol%)– Y^{3+} (0–50 mol%) codoped SiO_2 samples were annealed at 950 °C, all of them have an amorphous structure and all exhibit similar broad PL spectra (see Fig. 6). However, there still existed an apparent difference in FWHM between these spectra. The FWHM is ~45 nm for the Er³⁺ (1 mol%)–Y³⁺ (10 mol%) codoped SiO₂, which is larger than that of the Er³⁺ (1 mol%)-doped SiO₂ without the Y³⁺ codopant (FWHM = \sim 31 nm). The broadening of the PL spectra indicates that the Y³⁺ codopant plays a modifier role in affecting the bonding environment of Er³⁺ ions that can lead to a wider diversity of Er³⁺ bonding sites. However, for Er³⁺ (1 mol%)-Y³⁺ codoped SiO₂ samples with a Y³⁺ concentration varying form 10 to 50 mol%, the spectral bandwidth does not show an apparent difference. This implies that 10 mol% Y³⁺ codoping is enough to modify the Er bonding environment and maximize the Er site diversity in the SiO₂ matrix. The abovementioned phenomenon is also observed for the Er3+-Y3+ codoped SiO₂ samples with a larger amount of Er³⁺ (5– 10 mol%). The full width at half maximum (\sim 57 nm) is also larger than that (\sim 40 nm) of the Er³⁺ (5–10 mol%)-doped SiO₂ without Y^{3+} codoping.

The phase evolution of $Er^{3+}-Y^{3+}$ codoped SiO_2 samples is strongly dependent on the $Er^{3+}+Y^{3+}$ concentration and annealing temperature. As shown in Fig. 3, for the samples with an $Er^{3+}+Y^{3+}$ concentration greater than 30 mol% and annealed above 985 °C, the highly crystalline $Er_xY_{2-x}Si_2O_7$ phase forms in the host matrix. The TEM micrograph of an Er^{3+} (5 mol%)– Y^{3+} (50 mol%) codoped SiO_2 sample annealed at 1000 °C shows that many dark small droplet precipitates were observed (Fig. 11). $^{25-27}$ The energy-dispersive X-ray (EDX) analysis reveals that these precipitates contain Er and Y elements. As compared with XRD patterns (Fig. 2(a)), these precipitates should be the $Er_xY_{2-x}Si_2O_7$ phase. This indicates that most of the Er^{3+} ions are located on the well-defined sites of the $Er_xY_{2-x}Si_2O_7$ phase. Therefore, a number of sharp PL lines were observed as shown in Fig. 7.

Since both the content of hydroxy groups (i.e., quenching

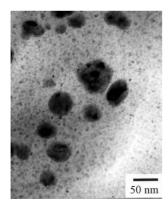


Fig. 11 TEM micrograph of Er^{3+} (5 mol%)– Y^{3+} (50 mol%) codoped SiO2 annealed at 1000 $^{\circ}C$ for 1 h.

centers) and the symmetry of local structure around the ${\rm Er}^{3+}$ ions can affect the PL intensity, the final-revealed PL intensity results from the competition between these two factors. For the samples annealed at 970 °C (*i.e.*, no crystalline ${\rm Er}_x{\rm Y}_{2-x}{\rm Si}_2{\rm O}_7$ phase existing in the host matrix), the varying PL intensity is primarily related to the amount of OH $^-$ hydroxy impurities. Therefore, in the temperature range 800–970 °C, the PL enhancement with increasing annealing temperature is mostly attributed to the decrease of the hydroxy quantity. $^{28-31}$

When the annealing temperature reached 985 °C, the poorly crystalline $\text{Er}_x Y_{2-x} \text{Si}_2 O_7$ phase had formed in the host matrix and resulted in the reduction of the probability of the $^4I_{13/2} \rightarrow ^4I_{15/2}$ transitions. However, a maximum PL intensity was observed at this temperature (985 °C) that reveals that the greatly decreased content of hydroxy quenching centers (see Fig. 5) can still offset the influence of the higher local symmetry of the Er site. On the other hand, as the samples were annealed at higher temperatures above 1000 °C, the $\text{Er}_x Y_{2-x} \text{Si}_2 O_7$ phase has been highly crystallized and thus the local structure around the Er^{3+} ions becomes more symmetric. Therefore, an abrupt reduction of the PL intensity is observed.

4.2. Role of Y^{3+} codopant on the development of \mbox{Er}^{3+} local structure

EXAFS analysis in Fig. 10 shows that there is local Er₂O₃-like structure in the amorphous Er3+-Y3+ codoped SiO2 host matrix. These Er sites with Er₂O₃-like local symmetry are generally thought of as the active luminescent centers. 32-34 Because Er³⁺ and Y³⁺ have the same valence and similar ionic radii (0.0881 and 0.0892 nm, respectively), they could be replaced by each other. Therefore, by codoping Y³⁺ ions into the Er³⁺-doped SiO₂ network, we believe that the –Er–O–Er– O-Er-bonding structure can be possibly changed into -Er-O- $(Y-O)_n$ -Er-, which indicates that the average interionic distance between Er^{3+} ions can be enlarged. Additionally, the FTIR spectra (see Fig. 4) shows that the addition of a large number of Y³⁺ codopant could destroy the network of SiO₂, leading to an increase of non-bridging oxygen groups in the SiO_2 matrix as observed in the Er^{3+} -doped SiO_2 system by Al^{3+} codoping. Some works reported that the non-bridging oxygen groups can reduce the tendency of the Er³⁺ ions to cluster. ^{10,13} According to the above-mentioned mechanism, the Y³⁺ codopant plays an important role in increasing the dispersion and solubility of Er³⁺ ions in the amorphous Er³⁺ Y³⁺ codoped SiO₂ systems, which results in a reduction of the concentration quenching effect and an increase of PL intensity.

These explanations are very consistent with our experimental results evidenced in Fig. 9. When the Y^{3+} codoping concentration is 0 and 10 mol%, the variation of PL intensity shows the phenomenon: $I_{(1\%)} \cdot I_{(5\%)} \cdot I_{(10\%)}$ [where $I_{(1\%)}$ represents the PL intensity of the sample with 1 mol% Er^{3+} doping dose]. However, for the 20 mol% Y^{3+} codoping concentration, the

phenomenon changes to $I_{(5\%)} \cdot I_{(10\%)} \cdot I_{(1\%)}$. This indicates that $10 \text{ mol}\% \ \text{Y}^{3+}$ codoping concentration is still not enough to disperse Er^{3+} ions very well and a large amount of Er^{3+} ions (5 and 10 mol%) still have access to form clusters. When a $20 \text{ mol}\% \ \text{Y}^{3+}$ codoping concentration was used, however, the influence of the concentration quenching effect on $I_{(5\%)}$ and $I_{(10\%)}$ can be considerably reduced.

For the sample with Y^{3+} concentration above 30 mol%, as the poorly crystalline pyrosilicate phase $(Er_xY_{2-x}Si_2O_7)$ was crystallized in the host matrix, it can be assumed that the Er^{3+} ions are located in the $Y_2Si_2O_7$ matrix. If the Er^{3+} ions were postulated to randomly disperse in the $Er_xY_{2-x}Si_2O_7$ phases, then the average spatial distance between Er^{3+} ions should be enlarged because some Er sites were occupied by $Er}^{3+}$ ions. This indicates that the concentration quenching effect can be reduced and hence the PL intensity is enhanced for the Er^{3+} (5–10 mol%)– $Er}^{3+}$ (30–50 mol%) codoped $Er}^{3+}$ (30–50 mol%) codopant still efficiently disperses the Er^{3+} (5–10 mol%) ions, which can offset the PL intensity loss resulting from the symmetry effect of the Er site.

5. Conclusion

Er³+–Y³+ codoped SiO₂ powdered bulks were prepared by a sol–gel process. The maximum ~1.54 μm PL intensity was obtained for the Er³+ (10 mol%)–Y³+ (50 mol%) codoped SiO₂ sample annealed at 985 °C. This can be attributed to the competition between the content of hydroxy groups and Er site symmetry. Below 985 °C, the content of hydroxy groups plays an important role in PL intensity. On the other hand, above 1000 °C, the highly crystalline Er_xY₂–_xSi₂Oγ phase forms and the local environment around the Er³+ ions becomes more symmetrical, resulting in reduced PL intensity and better resolved PL spectra. Additionally, the Y³+ codopant not only affects the crystallization behavior of the Er³+–Y³+ codoped SiO₂ sample but also modifies the bonding environment of Er³+ ions, which leads to an enlarged interionic distance between two nearby Er³+ ions and a wider diversity of Er³+ bonding sites. Therefore, a larger bandwidth of ~1.54 μm PL spectrum with improved efficiency is obtained for the Er³+–Y³+ codoped SiO₂ system.

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References

- 1 M. J. Weber, J. Non-Cryst. Solids, 1990, 123, 208.
- 2 B. Pedersen, W. J. Miniscalco and R. S. Quimby, *IEEE Photonics Technol. Lett.*, 1992, 4, 446.
- 3 R. Reisfeld and Y. Eckstein, J. Chem. Phys., 1975, 63, 4001.
- 4 W. F. Krupke and J. B. Gruber, J. Chem. Phys., 1963, 39, 1024.
- 5 W. F. Krupke and J. B. Gruber, *J. Chem. Phys.*, 1964, **41**, 1225.
- 6 E. Bescher and J. D. Mackenzie, *Mater. Sci. Eng. C*, 1998, 6(2–3), 145.
- 7 K. Kojima, K. Tsuchiya and N. Wada, J. Sol–Gel Sci. Tech., 2000, 19 511
- 8 B. R. Judd, Phys. Rev., 1962, 127, 750.
- R. M. Moon, W. C. Koehler, H. R. Child and L. J. Raubenheimer, *Phys. Rev.*, 1968, 176, 722.
- K. Arai, H. Namikawa, K. Kumata, T. Honda, Y. Ishii and T. Handa, J. Appl. Phys., 1986, 59, 3430.
- 11 C. K. Ryu, H. Choi and K. Kim, Appl. Phys. Lett., 1995, 66, 2496.
- 12 B. J. Ainslie, S. P. Craig and S. T. Davey, *Mater. Lett.*, 1987, 5, 143.

- 13 Y. Zhou, Y. L. Lam, S. S. Wang, H. L. Liu, C. H. Kam and Y. C. Chan, Appl. Phys. Lett., 1997, 71, 587.
- 14 ASTM JCPDS International Centre for Diffraction Data, 1997 [Ref.: 24-0062: Pfoertsch].
- 15 ASTM JCPDS International Centre for Diffraction Data, 1997 [Ref.: 32-1448 and 40-0034: Pfoertsch].
- 16 A. M. Efimov, J. Non-Cryst. Solids, 1996, 203, 1.
- 17 L. G. Hwa, S. L. Hwang and L. C. Liu, J. Non-Cryst. Solids, 1998, 238, 193.
- 18 A. A. Salem, R. Kellner and M. Grasserbauer, Glass Technol., 1994, 35, 135.
- 19 N. A. El-Shafi and M. M. Morsi, J. Mater. Sci., 1997, 32, 5185.
- 20 T. Ishikawa and S. Akagi, Phys. Chem. Glasses, 1987, 19, 108.
- K. Sun, W. H. Lee and W. M. Risen, Jr, J. Non-Cryst. Solids, 1987, 92, 145.
- 22 R. F. Bartholomew, B. L. Butler, H. L. Hoover and C. K. Wa, J. Am. Ceram. Soc., 1980, 63, 481.
- 23 K. Hafidi, Y. Ijdiyaou, M. Azizan, E. L. Ameziane, A. Outzourhit, T. A. Nguyen and M. Brunel, Appl. Surf. Sci., 1997, 108, 251.
- 24 Y. I. Smolin and Y. F. Shepelev, Acta Crystallogr. Sect. B: Struct. Crystallogr. Cryst. Chem., 1970, 26, 484.

- 25 L. L. Lee and D. S. Tsai, J. Mater. Sci. Lett., 1994, 13, 615.
- 26 W. V. Moreshead, J. R. Noguès and R. H. Krabill, *J. Non-Cryst. Solids*, 1990, **121**, 267.
- 27 A. Polman, J. Appl. Phys., 1997, 82, 1.
- 28 Y. Yan, A. J. Faber and H. de Waal, J. Non-Cryst. Solids, 1995, 181, 283.
- 29 E. Snoeks, P. G. Kik and A. Polman, Opt. Mater., 1996, 5, 159.
- M. Benatsou, B. Capoen, M. Bouazaoui, W. Tchana and J. P. Vilcot, Appl. Phys. Lett., 1997, 71, 428.
- 31 J. Phalippou, T. Woignier and J. Zarzycki, in *Ultrastructure Processing of Ceramics, Glasses and Composites*, Eds. L. L. Hench and D. R. Ulrich, Wiley, New York, 1984, p. 70.
- 32 D. L. Adler, D. C. Jacobson, D. J. Eaglesham, M. A. Marcus, J. L. Benton, J. M. Poate and P. H. Citrin, *Appl. Phys. Lett.*, 1992, 61, 2181.
- 33 C. Piamonteze, A. C. Iñiguez, L. R. Tessler, M. C. Martins Alves and H. Tolentino, *Phys. Rev. Lett.*, 1998, **81**, 4652.
- 34 A. Terrasi, G. Franzò, S. Coffa, F. Priolo, F. D'Acapito and S. Mobilio, J. Appl. Phys., 1997, 70, 1712.